

Geometric quantization of Lagrangian torus fibrations

and: a fact about integral-integral affine geometry

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Motivation: "invariance of polarization" phenomenon
in geometric quantization.

affine maps

$$\mathbb{R}^n \longrightarrow \mathbb{R}^n$$
$$x \longmapsto Ax + b$$

integral affine:

integral-integral affine:

A	b
$GL(n, \mathbb{R})$	\mathbb{R}^n
$GL(n, \mathbb{Z})$	\mathbb{R}^n
$GL(n, \mathbb{Z})$	\mathbb{Z}^n

B integral-integral affine mfd:

\Leftrightarrow mfd + maximal integral-integral affine atlas.

transition maps are locally integral-integral affine

\Rightarrow

• flat torsion free connection on TB

• integral lattice $\Lambda \subset TB$

period lattice $\Lambda^* \subset T^*B$

\Rightarrow measure on B

• integral points $B_{\mathbb{Z}} \subset B$

Fact:

B cpct $\Rightarrow |B_{\mathbb{Z}}| = \text{vol}(B)$

Examples: $B = \mathbb{R}^n / \Gamma$, $\Gamma \subset \mathbb{R}^n$

B	n	Γ generated by...
Torus	$n \geq 1$	$x_j \mapsto x_j + 1$ for $j=1, \dots, n$
Klein bottle	$n=2$	$x_1 \mapsto x_1 + 1$, $(x_1, x_2) \mapsto (-x_1, x_2 + 1)$
Kodaira-Thurston-like	$n \geq 3$	$(x_1, x_2, x_3) \mapsto (x_1 + x_2, x_2, x_3 + 1)$ $x_j \mapsto x_j + 1$ for $j \neq 3$

Question:

B integral integral affine, cpct $\stackrel{?}{\implies} B \cong \mathbb{R}^n / \Gamma$

Proper Lagrangian fibration :

(M, ω)

$\pi \downarrow$

B

Arnold - Liouville charts :

semilocal coordinates:

$M \supset \pi^{-1}(u)$

$(x_1, \dots, x_n, t_1, \dots, t_n)$

$t_j \text{ mod } 1$

\downarrow

\downarrow

$\omega = \sum_{j=1}^n dx_j \wedge dt_j$

$B \supset u$

(x_1, \dots, x_n)

\Rightarrow integral affine str on B

Prequantize :

$$(L, \langle \rangle, \nabla)$$

curvature $\nabla = \omega$



$$(M, \omega)$$



$$B$$

Enhanced Arnold Liouville charts :

Arnold-Liouville charts + local trivializations of L

s.t.
$$\nabla = d + i \sum_j x_j dt_j$$

\Rightarrow integral-integral affine str. on B

$$M_0 := T^*B/\Lambda^*$$

Lagrangian torus fibration



$$M^\vee := T B/\Lambda$$

semi-local coordinates :



$$\underbrace{x_1, \dots, x_n}_{\text{integral-integral local coordinates on } B}, \quad \underbrace{y_1, \dots, y_n}_{\text{mod } 1 \text{ adapted fibre coordinates}}$$

integral-integral
local coordinates
on B

mod 1
adapted fibre coordinates

Local torus \mathbb{T}^n actions on M^\vee
up to $\text{Aut}(\mathbb{T}^n)$ on chart overlaps

$$\alpha \in \Omega^k(W), \quad W = \pi^{-1}(U) \underset{\text{open}}{\subset} M^v$$

" α is \mathbb{T}^n -invariant"

$$\Leftrightarrow \text{locally } \alpha = g^* \alpha \quad \forall g \in \mathbb{T}^n$$

is well defined.

$$\Rightarrow \text{Complex } (\Omega_{\text{invt}}^*(M^v), d)$$

$$\Rightarrow \text{Cohomology } H_{\text{invt}}^*(M^v)$$

$$\Omega_{\text{invt}}^*(M^v) \xrightarrow{i} \Omega^*(M^v)$$

$$\Rightarrow i_*: H_{\text{invt}}^*(M^v) \longrightarrow H^*(M^v)$$

Lemma i_* is an isomorphism.

pf for global \mathbb{T}^n action:

$$\text{averaging } \bar{\alpha} := \int_{g \in \mathbb{T}^n} (g^* \alpha) \, dg$$

$$d\bar{\alpha} = \overline{d\alpha}.$$

$$\mathbb{T}^n \text{ connected} \Rightarrow \int_{\mathbb{T}^n} [g^* \alpha] = [\alpha] \Rightarrow [\bar{\alpha}] = [\alpha]$$

$\therefore H_{\text{int}}^i \rightarrow H^i$ is onto.

if α is invariant and $\alpha = d\beta$.

Then also $\alpha = d\bar{\beta}$.

$\therefore H_{\text{int}}^i \rightarrow H^i$ is one-to-one.

pf for local \mathbb{T}^n -action: induction on size of covering
by open sets w/ \mathbb{T}^n actions.

$l=1$: Case of global T^n action.

$l \rightarrow l+1$: Mayer Vietoris + five lemma:

$$\dots \rightarrow H_{inv}^k \left(\bigcup_{i=1}^{l+1} W_i \right) \rightarrow H_{inv}^k \left(\bigcup_{i=1}^l W_i \right) \oplus H_{inv}^k (W_{l+1}) \rightarrow H_{inv}^k \left(\left(\bigcup_{i=1}^l W_i \right) \cap W_{l+1} \right) \rightarrow \dots$$

$$\downarrow i_*$$

$$\downarrow i_* \oplus i_*$$

$$\downarrow i_*$$

$$\dots \rightarrow H^k \left(\bigcup_{i=1}^{l+1} W_i \right) \rightarrow H^k \left(\bigcup_{i=1}^l W_i \right) \oplus H^k (W_{l+1}) \rightarrow H^k \left(\left(\bigcup_{i=1}^l W_i \right) \cap W_{l+1} \right) \rightarrow \dots$$



Note. $av_* \circ i_* = \text{Identity on } H_{inv}^*$

$\&$ i_* is an isomorphism

\Rightarrow also $i_* \circ av_* = \text{Identity on } H^*$

B integral affine, $\Lambda \subset TB$ integral lattice.

$$M^v = TB/\Lambda$$

\downarrow
 B

$$x_1, \dots, x_n, y_1, \dots, y_n$$

$\underbrace{\hspace{10em}}_{\text{mod } \mathbb{Z}}$

semi-local
adapted
coordinates

$$Z_0 := \{y_j = 0\} \quad \text{zero section.}$$

Assume B is oriented, & x_1, \dots, x_n are compatible with the orientation.

Then.

- $dy_1 \wedge \dots \wedge dy_n$ is independent of the choice of adapted coordinates
- $[dy_1 \wedge \dots \wedge dy_n] = \text{P.D.}(Z_0)$

pf.

Take any coh. class $[\alpha] \in H^n$.

wlog α is invt.

Locally $\alpha = a(x) dx_1 \wedge \dots \wedge dx_n$.
+ terms with dy_j 's

$$\int_{Z_0} \alpha = \int_{Z_0} a(x) dx_1 \wedge \dots \wedge dx_n$$

$$\stackrel{\text{Fubini}}{=} \int_{M^v} a(x) dx_1 \wedge \dots \wedge dx_n \wedge dy_1 \wedge \dots \wedge dy_n$$

$$= \int_{M^v} \alpha \wedge (dy_1 \wedge \dots \wedge dy_n)$$

$Z_0 := \{y_j = 0\} \subset M^v$ zero section

$D := \{y_j = x_j \pmod{1}\} \subset M^v$ diagonal section.

$$|Z_0 \cap D| = |B_Z|$$

pf that $|B_Z| = \text{vol}(B)$:

$$|B_Z| = |D \cap Z_0| = \left| \int_D [dy_1 \wedge \dots \wedge dy_n] \right|$$

$$= \left| \int_B \underbrace{\omega^* dy_1 \wedge \dots \wedge dy_n}_{= dx_1 \wedge \dots \wedge dx_n} \right|$$

$$\sigma: x_j \mapsto (x_j, x_j \pmod{1})$$

$$= \text{vol}(B).$$

Bohr-Sommerfeld quantization



$x \in B$ Bohr-Sommerfeld
 $\Leftrightarrow (L|_{\pi^{-1}(x)}, \nabla)$ is trivial

$\Leftrightarrow x \in B_{\text{reg}}$.

The Hilbert space: $\mathcal{H}_{B_S} := \bigoplus_{\text{Bohr-Sommerfeld}} \mathbb{C}$.

Dirac-Dolbeault quantization:

The Hilbert space: \mathcal{H} s.t.

$$\dim \mathcal{H} = RR(M, \omega) := \int_M e^\omega \text{Todd}(TM, \mathcal{J}).$$

The RHS:

cptble a.c. str.

$$e^\omega = 1 + \omega + \frac{\omega^2}{2!} + \dots$$

$$\int_{M^v} e^\omega = \int_M \frac{\omega^n}{n!}$$

Liouville volume

$$\text{Todd}(x) := \prod_j \frac{x_j}{1 - e^{-x_j}}$$

$$\rightsquigarrow \text{Todd}: \text{Lie}(\mathfrak{u}(n)) \rightarrow \mathbb{R}$$

$$\rightsquigarrow \text{Chern Weil} \quad \text{Todd}: \left\{ \begin{array}{l} \text{cplx v. bundles} \\ \text{over } M \end{array} \right\} \longrightarrow H^*(M, \mathbb{R})$$

The LHS: e.g., can take

$$H = \Gamma_{\text{hol}}(L)$$

or $H = \sum (-1)^i H^i(\mathcal{O}_L)$

or $H = \ker \bar{\partial}_L - \text{coker } \bar{\partial}_L$

Theorem:

$$RR(\mu, \omega) = |B_{\mathbb{Z}}|$$

Bohn
Sommerfeld
set.

interpretation: "invariance of polarization"

Proof

$$V := \ker \pi_* \subset TM$$

$$\begin{array}{c} \mathbb{Z} \\ \downarrow \\ M \\ \downarrow \\ B \end{array}$$

$$0 \rightarrow V \rightarrow TM \rightarrow \pi^* TB \rightarrow 0$$

vector bundles over M

$$V \cong \pi^* T^*B$$

$$V \otimes \mathbb{C} \cong (TM, \mathcal{J})$$

$$\text{via } u + iv \mapsto u + \mathcal{J}v$$

$$u, v \in V_y \quad y \in M$$

$$\Rightarrow (TM, \mathcal{J}) \cong \pi^* \underbrace{(T^*B \otimes \mathbb{C})}_{\text{flat}}$$

$\Rightarrow (TM, J)$ is flat.

$$\Rightarrow \text{Todd}(TM, J) = 1$$

$$\Rightarrow \text{RR}(\overset{M}{\omega}, \omega) = \int_M e^{\omega} \underbrace{1}_{\text{Todd}}$$

= Liouville volume of M

Fubini: = integral of affine volume of B

$$= |B_{\mathbb{Z}}|$$

By the fact
from integral-integral
geometry