Rigidity phenomena in Strongly correlated random point fields and The emergence of forbidden regions

> Subhro Ghosh National University of Singapore

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Poisson Process

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Poisson Process Ginibre Eigenvalues

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Process	Eigenvalues	Zeroes

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Subhro Ghosh National University of Singapore Rigidity Phenomena

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- $\bullet\,$ In Poisson point process, the points inside and outside ${\cal D}$ are independent of each other

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• Finite *n*: μ_n = Eigenvalues of $G_n = ((\xi_{ij}))_{1 \le i,j \le n}$, ξ_{ij} i.i.d $N_{\mathbb{C}}(0,1)$ (NO normalization by \sqrt{n})

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• Finite n:
$$f_n(z) = \xi_0 + \frac{\xi_1}{\sqrt{1!}}z + \dots + \frac{\xi_k}{\sqrt{k!}}z^k + \dots + \frac{\xi_n}{\sqrt{n!}}z^n$$

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- Arises in the study of quantum chaotic dynamics, eg. in the work of Bogomolnyi, Bohigas, Leboeuf.

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 $\Sigma_{S(\omega)}$: constant sum hypersurface $\sum_{i=1}^{N(\omega)} \zeta_i = S(\omega)$ inside $\mathcal{D}^{N(\omega)}$

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- π is rigid at level k if
 - The points of π outside \mathcal{D} determine $0, 1, \ldots, (k-1)$ moments of the points in \mathcal{D}

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- Natural point processes for Levels $k \ge 3$??
A Hierarchy of Rigidity phenomena

Theorem (G.-Krishnapur, Comm. Math. Phys. '21)

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• Phase transition in the rigidity behaviour in α at the values $\alpha = 1, \frac{1}{2}, \frac{1}{3}, \frac{1}{4}, \cdots$ (rigid at level k for $\alpha \in (\frac{1}{k}, \frac{1}{k-1}]$).

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- The standard planar GAF is Gaussian r.v. in the space of L^2 analytic functions w.r.t. complex Gaussian measure. The α -GAFs are Gaussian r.v. in the space of L^2 analytic functions w.r.t. other measures on \mathbb{C} (roughly, density $\sim e^{-|z|^{2/\alpha}}$).

Theorem (Gangopadhyay-G.-Tan, Commun. Pure Appl. Math. '23)

For the general α -Gaussian zero ensembles, as well as the Ginibre ensemble, there are positive quantities $m(\omega)$, $M(\omega)$ such that the conditional density f_{ω} satisfies, on its support $\Xi(\omega)$, the following :

$$m(\omega) \exp(2\sum_{i\neq j} \log |\zeta_i - \zeta_j|) \mathbb{1}_{\Xi(\omega)} \leq f_{\omega}(\zeta),$$

and

$$f_{\omega}(\zeta) \leq M(\omega) \exp(2\sum_{i \neq j} \log |\zeta_i - \zeta_j|) \mathbb{1}_{\Xi(\omega)},$$

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Thus, almost surely the conditional measure exhibits quadratic repulsion. Provides a general programme for demonstrating approximate Gibbsian structures in strongly correlated point processes, which can be executed on the basis of finite particle estimates.

Theorem (Reconstruction of Gaussian Analytic Function, G.-Peres)

The zeroes of the GAF determine the function a.s. (up to a multiplicative factor of modulus 1). In other words, if ν denotes the zeroes of the GAF f, then \exists an analytic function $g(z) = \sum_{k=0}^{\infty} a_k(\nu) z^k \text{ such that } f(z) = \gamma . g(z)$ Here γ follows Unif(S¹) and is independent of ν .

Theorem (G., Elec. Comm. Probab. '16)

For a point process P on a background space $\Xi(\mathbb{R} \text{ or } \mathbb{C})$ which is rigid at level k, and point configurations $x \in \Xi^m, y \in \Xi^n$, the reduced Palm measures P_x and P_y are mutually absolutely continuous if and only if we have matching moments $m_i(x) = m_i(y)$ for $i \le k$.

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- In particular, for the GAF zero process, the Palm measures ρ_z and ρ_w are mutually singular for a.e. z, w (in spite of translation invariance).
- Rigidity phenomena have subsequently seen extensive works by Bufetov, Dereudre, Leble, Maida, Leble, Shirai, Najnudel

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- For Gaussian random polynomials of high degree (growing as αR^2)The probability of this event is $\simeq e^{-cR^4}$.
- The constant *c* is explicitly known.

Question (Key questions)

• Given that there is a hole of size R, how do the zeroes look like, even in expectation ?

 Comparison with GUE eigenvalue process on ℝ: the conditional intensity exhibits a spike at R and subsequently decreases to the equilibrium intensity. (Majumdar et. al.)

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- Comparison with Ginibre eigenvalue ensemble on C: the conditional intensity exhibits a singular component at R and subsequently decreases to the equilibrium intensity. (Lebowitz et. al., Shirai)

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- Beyond $\sqrt{e}R$, the conditional intensity behaves, in the limit $R \rightarrow \infty$, like the equilibrium intensity.

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- There is a singular component at the edge of the hole
- There is subsequent "forbidden region", namely, in the annulus $R < r < \sqrt{e}R$, the conditional intensity $\rightarrow 0$ as $R \rightarrow \infty$.

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- Beyond $\sqrt{e}R$, the conditional intensity behaves, in the limit $R \rightarrow \infty$, like the equilibrium intensity.
- Subsequent results on holes of more general shape (Nishry-Wenmann)



Gaussian Matrices

Subhro Ghosh National University of Singapore Rigidity Phenomena

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Matrices

Gaussian Polynomials

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• Upcoming work : Existence of *Forbidden Regions* on Riemann surfaces (Dinh-G.-Wu)

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- Know outside zeroes \Rightarrow Know $\int_{\mathcal{D}_L \setminus \mathcal{D}} \varphi_L d\nu \Rightarrow$ Compute $n(\mathcal{D})$ approximately, now let $L \to \infty$

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💐 💐 💐 Happy Birthday to Prof Dinh !! 💐 💐 💐

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- Forbidden regions for random zeros on Riemann surfaces, with T.C. Dinh and H. Wu (near completion)