

Conformal Loop Ensemble for the Integer Quantum Hall Transition

E. Bettelheim

Hebrew University, Jerusalem

Workshop on Geometric aspects of the Quantum Hall Effect
December 17, 2015
Cologne, Germany

Acknowledgement

Introduction

CLE Approach

Conclusion

Work with I. Gruzberg.

Acknowledgement

Introduction

-Chalker-Coddington

-Positive Weights

-Restriction

-Recap

CLE Approach

Conclusion

Introduction

The Chalker-Coddington Network

Acknowledgement

Introduction

-Chalker-Coddington

-Positive Weights

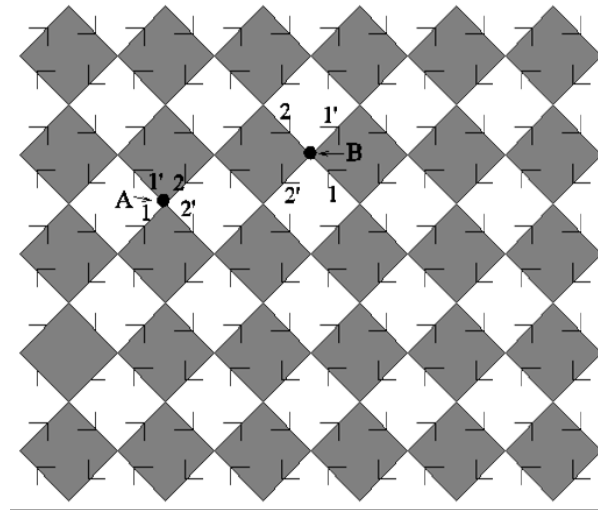
-Restriction

-Recap

CLE Approach

Conclusion

- The Chalker-Coddington network:



The Chalker-Coddington Network

Acknowledgement

Introduction

-Chalker-Coddington

-Positive Weights

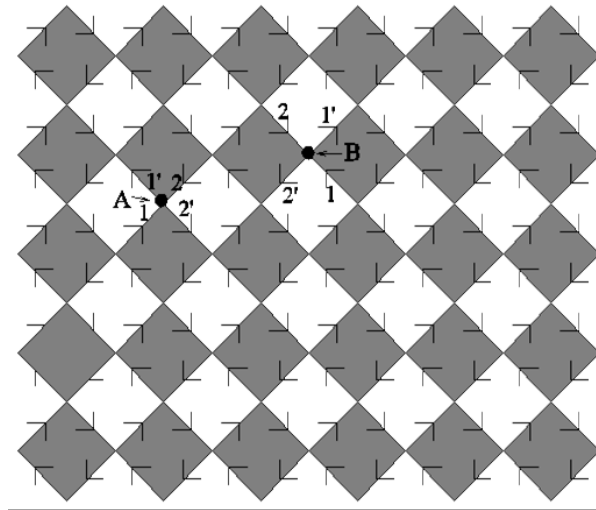
-Restriction

-Recap

CLE Approach

Conclusion

- The Chalker-Coddington network:



- The wavefunction accumulates a random phase on each link

The Chalker-Coddington Network

Acknowledgement

Introduction

-Chalker-Coddington

-Positive Weights

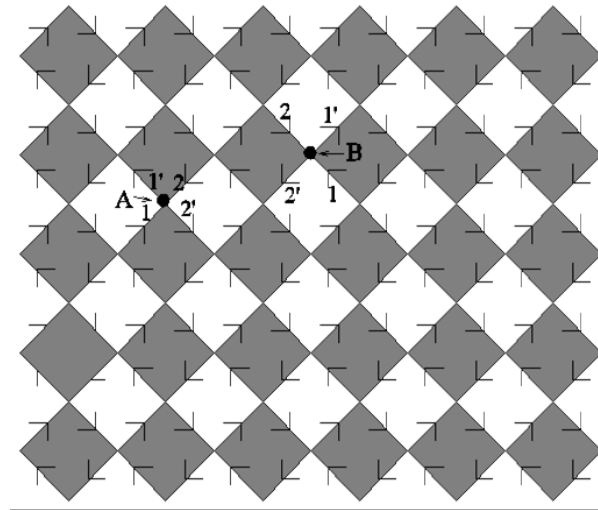
-Restriction

-Recap

CLE Approach

Conclusion

- The Chalker-Coddington network:



- The wavefunction accumulates a random phase on each link
- A beam splitter is at each node:

$$S = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix}$$

The Chalker-Coddington Network

Acknowledgement

Introduction

-Chalker-Coddington

-Positive Weights

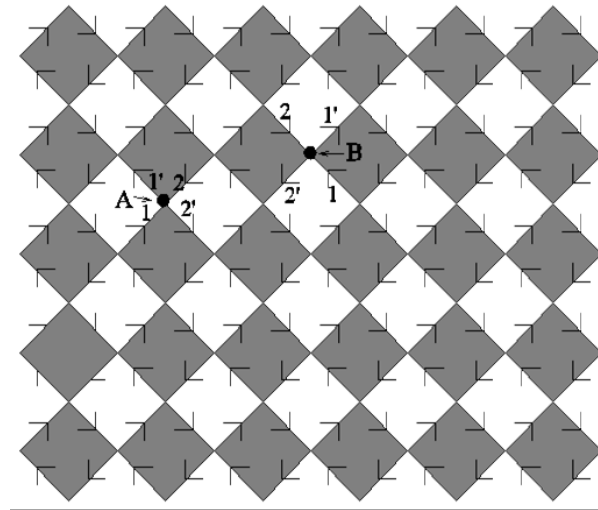
-Restriction

-Recap

CLE Approach

Conclusion

- The Chalker-Coddington network:



- The probability amplitude from a to b is:

$$G(a, b) = \sum_P \prod_{l \in P} e^{i\theta_l} \prod_{(l, l') \in P} S_{l, l'}$$

$$S = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix}$$

Positive Weights

Acknowledgement

Introduction

-Chalker-Coddington

-Positive Weights

-Restriction

-Recap

CLE Approach

Conclusion

- We compute

$$\sigma(a, b) = \langle |G(a, b)|^2 \rangle = \sum_{P, P'} \langle W(P) W^*(P') \rangle,$$

with $W(P) = \prod_{l \in P} e^{i\theta_l} \prod_{(l, l') \in P} S_{l, l'}$.

Positive Weights

Acknowledgement

Introduction

-Chalker-Coddington

-Positive Weights

-Restriction

-Recap

CLE Approach

Conclusion

- We compute
$$\sigma(a, b) = \langle |G(a, b)|^2 \rangle = \sum_{P, P'} \langle W(P) W^*(P') \rangle,$$
with $W(P) = \prod_{l \in P} e^{i\theta_l} \prod_{(l, l') \in P} S_{l, l'}$.
- The phases need to cancel

Positive Weights

Acknowledgement

Introduction

-Chalker-Coddington

-Positive Weights

-Restriction

-Recap

CLE Approach

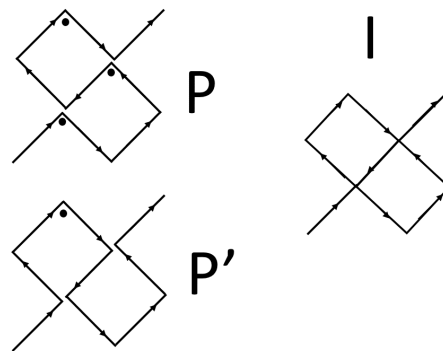
Conclusion

- We compute

$$\sigma(a, b) = \langle |G(a, b)|^2 \rangle = \sum_{P, P'} \langle W(P) W^*(P') \rangle,$$

with $W(P) = \prod_{l \in P} e^{i\theta_l} \prod_{(l, l') \in P} S_{l, l'}$.

- The phases need to cancel \implies the picture of P is the same as P'



Positive Weights

Acknowledgement

Introduction

-Chalker-Coddington

-Positive Weights

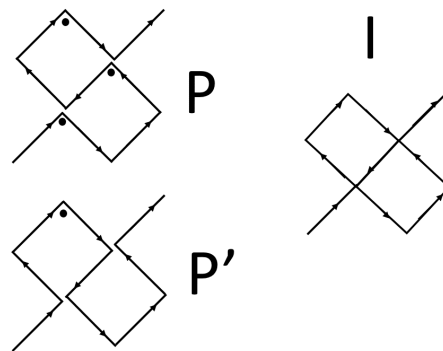
-Restriction

-Recap

CLE Approach

Conclusion

- We compute $\sigma(a, b) = \langle |G(a, b)|^2 \rangle = \sum_{P, P'} \langle W(P)W^*(P') \rangle$, with $W(P) = \prod_{l \in P} e^{i\theta_l} \prod_{(l, l') \in P} S_{l, l'}$.
- The phases need to cancel \implies the picture of P is the same as P'



- We get $(s = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix})$

$$\sigma(a, b) = \sum_I \frac{A(I)^2}{2^{L(I)}}, \quad A(I) = \sum_{P \in I} \prod_{(l, l') \in P} (-)^{\sigma(l, l')}$$

Conformal Restriction

Acknowledgement

Introduction

-Chalker-Coddington

-Positive Weights

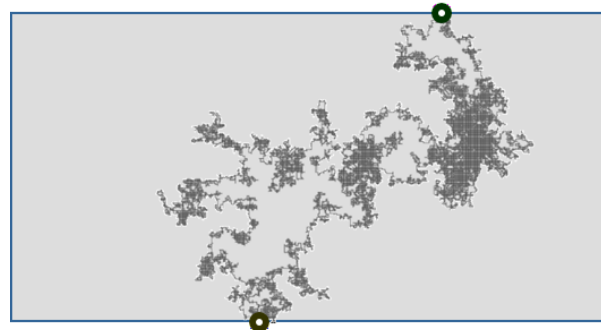
-Restriction

-Recap

CLE Approach

Conclusion

- A Large Cluster's Shape is determined EB, I. Gruzberg, A. W. W. Ludwig by conformal invariance and intrinsic weights ('conformal restriction' Lawler)



Conformal Restriction

Acknowledgement

Introduction

-Chalker-Coddington

-Positive Weights

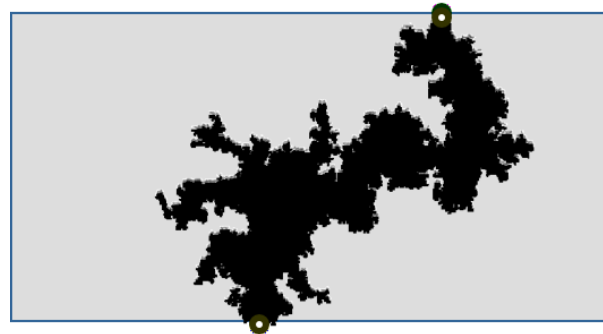
-Restriction

-Recap

CLE Approach

Conclusion

- A Large Cluster's Shape is determined EB, I. Gruzberg, A. W. W. Ludwig by conformal invariance and intrinsic weights ('conformal restriction' Lawler)



- Only external perimeter is determined.

Conformal Restriction

Acknowledgement

Introduction

-Chalker-Coddington

-Positive Weights

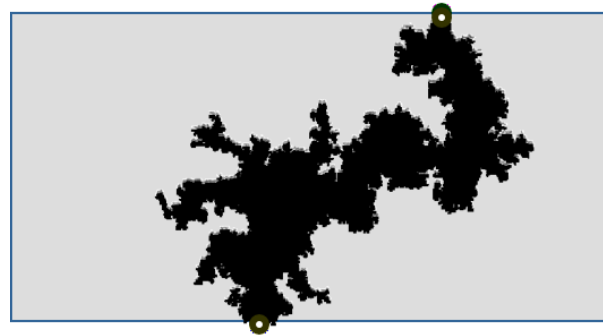
-Restriction

-Recap

CLE Approach

Conclusion

- A Large Cluster's Shape is determined EB, I. Gruzberg, A. W. W. Ludwig by conformal invariance and intrinsic weights ('conformal restriction' Lawler)



- Only external perimeter is determined.
- It's the same perimeter for Percolation, Random Walks, Transition, etc.

Recap

Acknowledgement

Introduction

-Chalker-Coddington

-Positive Weights

-Restriction

-Recap

CLE Approach

Conclusion

- A statistical approach may be obtained for the integer quantum Hall transition.

Recap

Acknowledgement

Introduction

-Chalker-Coddington

-Positive Weights

-Restriction

-Recap

CLE Approach

Conclusion

- A statistical approach may be obtained for the integer quantum Hall transition.
- The approach gives fractal properties of the electron's path.

Recap

Acknowledgement

Introduction

-Chalker-Coddington

-Positive Weights

-Restriction

-Recap

CLE Approach

Conclusion

- A statistical approach may be obtained for the integer quantum Hall transition.
- The approach gives fractal properties of the electron's path.
- The accessible fractal properties seem too general.

Recap

Acknowledgement

Introduction

-Chalker-Coddington

-Positive Weights

-Restriction

-Recap

CLE Approach

Conclusion

- A statistical approach may be obtained for the integer quantum Hall transition.
- The approach gives fractal properties of the electron's path.
- The accessible fractal properties seem too general.
- A more detailed description of the interior of the path is needed.

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

-Fermions

-The Loops

-Re-weighting

Conclusion

Conformal Loop Ensemble Approach

Conformal Loop Ensembles

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

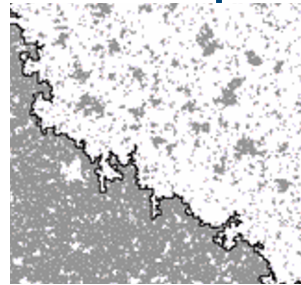
-Fermions

-The Loops

-Re-weighting

Conclusion

- In statistical mechanics, one can often describe the problem as that of loops in the plane.



Conformal Loop Ensembles

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

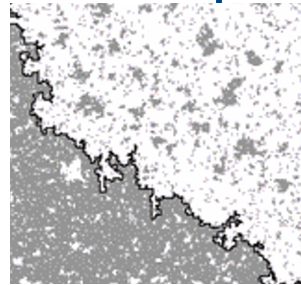
-Fermions

-The Loops

-Re-weighting

Conclusion

- In statistical mechanics, one can often describe the problem as that of loops in the plane.



- At critical points these loops have conformal invariance properties.

Conformal Loop Ensembles

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

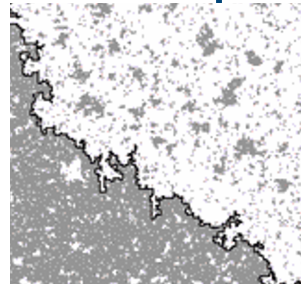
-Fermions

-The Loops

-Re-weighting

Conclusion

- In statistical mechanics, one can often describe the problem as that of loops in the plane.



- At critical points these loops have conformal invariance properties.
- The classification of possible critical models can be done either by appealing to either

Conformal Loop Ensembles

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

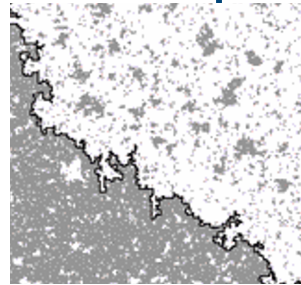
-Fermions

-The Loops

-Re-weighting

Conclusion

- In statistical mechanics, one can often describe the problem as that of loops in the plane.



- At critical points these loops have conformal invariance properties.
- The classification of possible critical models can be done either by appealing to either
 1. Conformal Field Theory

Conformal Loop Ensembles

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

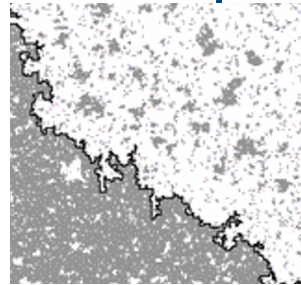
-Fermions

-The Loops

-Re-weighting

Conclusion

- In statistical mechanics, one can often describe the problem as that of loops in the plane.



- At critical points these loops have conformal invariance properties.
- The classification of possible critical models can be done either by appealing to either
 1. Conformal Field Theory or
 2. Loop Ensembles and the Schramm-Loewner Equation (SLE).

The Manhattan (medial) Lattice

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

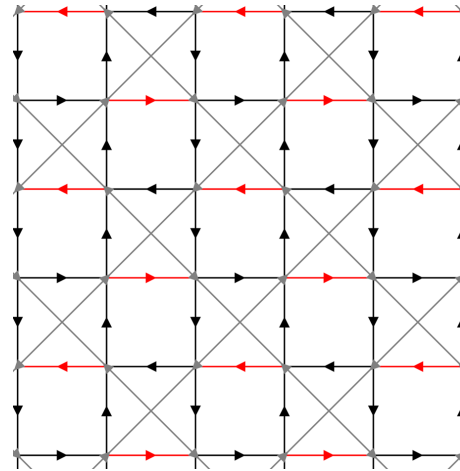
-Fermions

-The Loops

-Re-weighting

Conclusion

- It's easier to keep track of the sign on the median lattice, which in this case is the Manhattan lattice.



The Manhattan (medial) Lattice

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

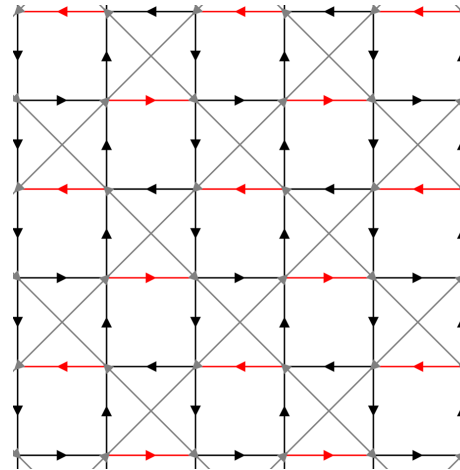
-Fermions

-The Loops

-Re-weighting

Conclusion

- It's easier to keep track of the sign on the median lattice, which in this case is the Manhattan lattice.



- We still need to count the ways to traverse the lattice, which interfere.

Trees

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

-Fermions

-The Loops

-Re-weighting

Conclusion

- Given a path, draw (for every node visited by the path) the link from which we got to that node first.

Trees

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

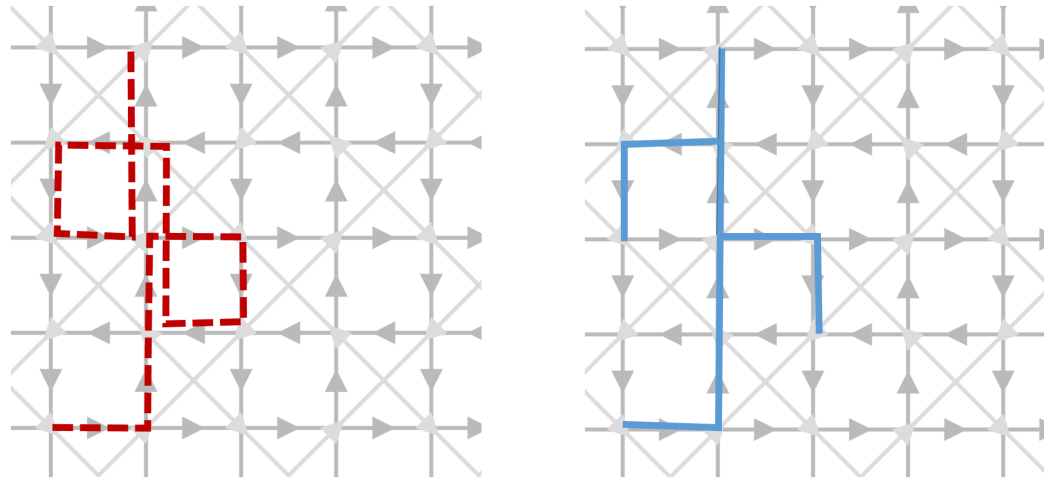
-Fermions

-The Loops

-Re-weighting

Conclusion

- Given a path, draw (for every node visited by the path) the link from which we got to that node first.



Trees

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

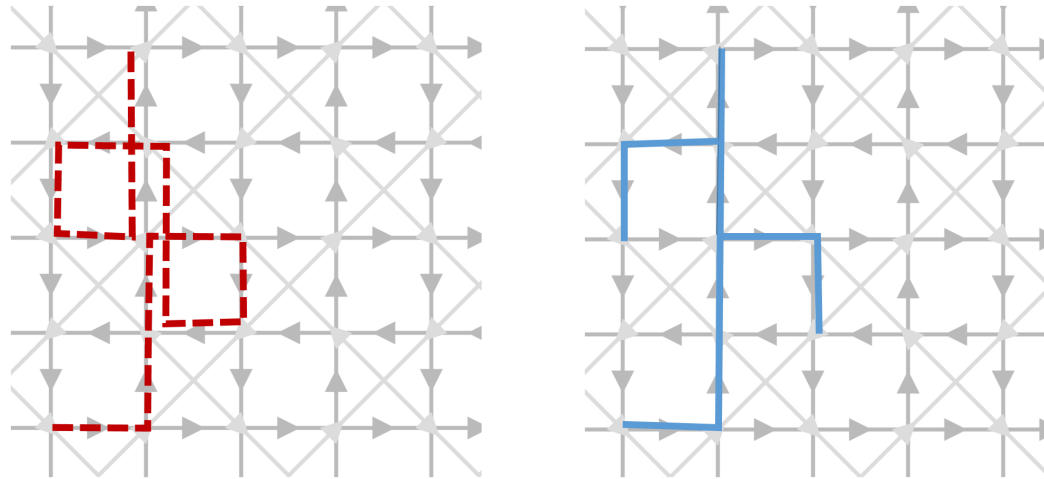
-Fermions

-The Loops

-Re-weighting

Conclusion

- Given a path, draw (for every node visited by the path) the link from which we got to that node first.



- One can reconstruct the possible paths given the tree and the number of visits to each link.

Trees

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

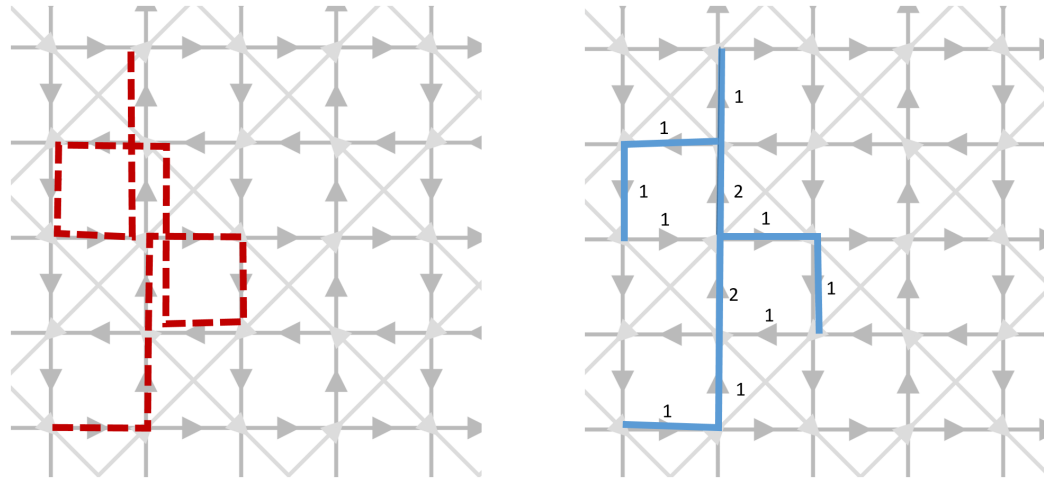
-Fermions

-The Loops

-Re-weighting

Conclusion

- Given a path, draw (for every node visited by the path) the link from which we got to that node first.



- One can reconstruct the possible paths given the tree and the number of visits to each link.

Free Fermions

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

-Fermions

-The Loops

-Re-weighting

Conclusion

- Dealing with tree is easier due to Kasteleyn, Temperley, Kenyon:

$$|T| = \det(K)$$

Free Fermions

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

-Fermions

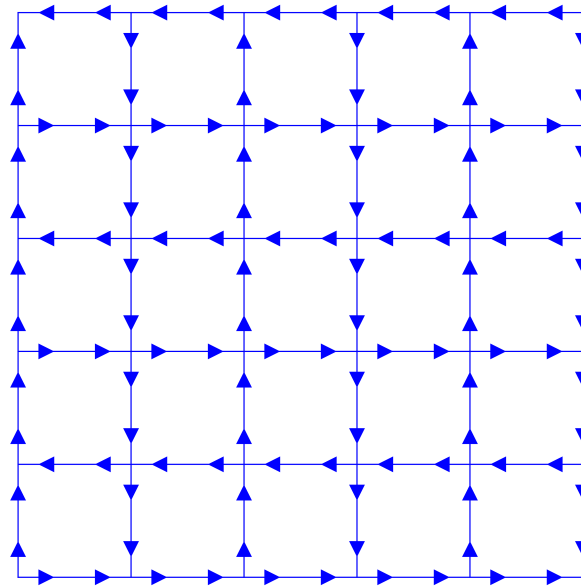
-The Loops

-Re-weighting

Conclusion

- Dealing with tree is easier due to Kasteleyn, Temperley, Kenyon:

$$|T| = \det(K)$$



Free Fermions

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

-Fermions

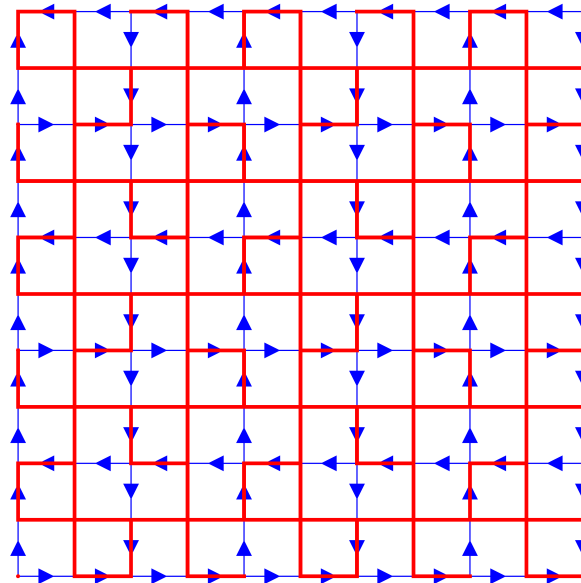
-The Loops

-Re-weighting

Conclusion

- Dealing with tree is easier due to Kasteleyn, Temperley, Kenyon:

$$|T| = \det(K)$$



Free Fermions

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

-Fermions

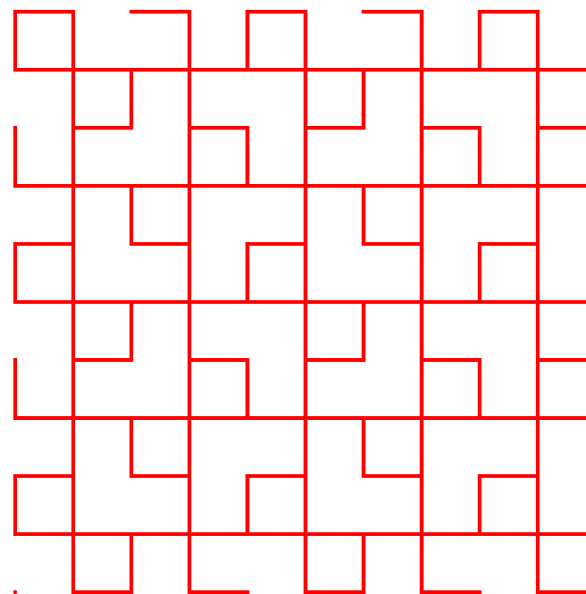
-The Loops

-Re-weighting

Conclusion

- Dealing with tree is easier due to Kasteleyn, Temperley, Kenyon:

$$|T| = \det(K)$$



Free Fermions

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

-Fermions

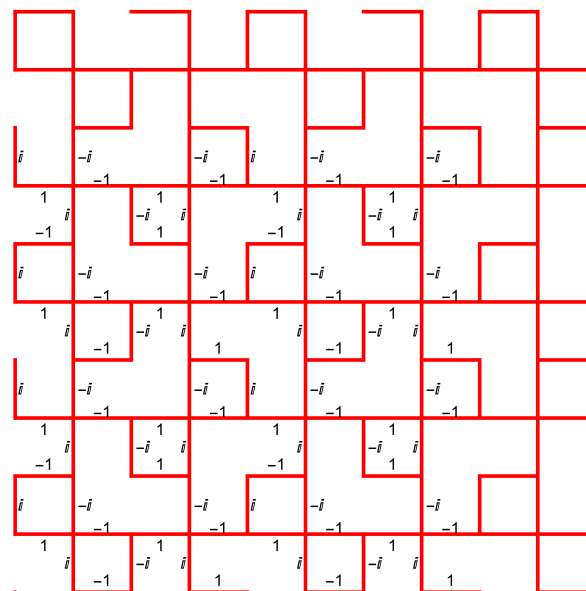
-The Loops

-Re-weighting

Conclusion

- Dealing with tree is easier due to Kasteleyn, Temperley, Kenyon:

$$|T| = \det(K)$$



Free Fermions

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

-Fermions

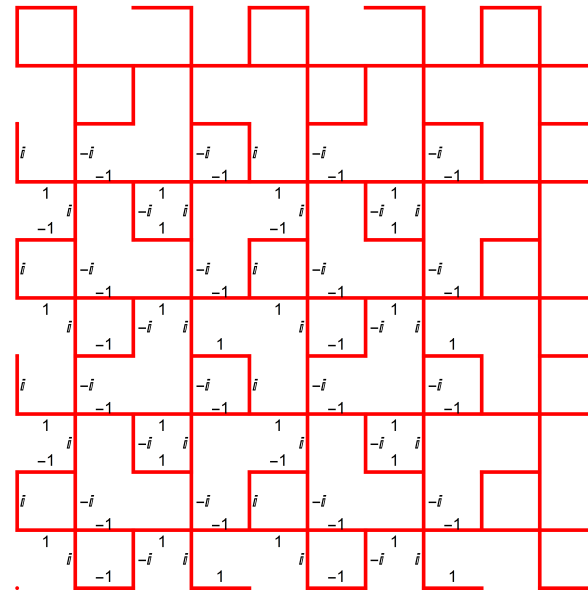
-The Loops

-Re-weighting

Conclusion

- Dealing with tree is easier due to Kasteleyn, Temperley, Kenyon:

$$|T| = \det(K)$$



- $\langle \psi_{i_1}^\dagger \psi_{j_1} \psi_{i_2}^\dagger \psi_{j_2} \dots \rangle = \det_{k,l} (K^{-1}(i_k, j_l))$

The Loops

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

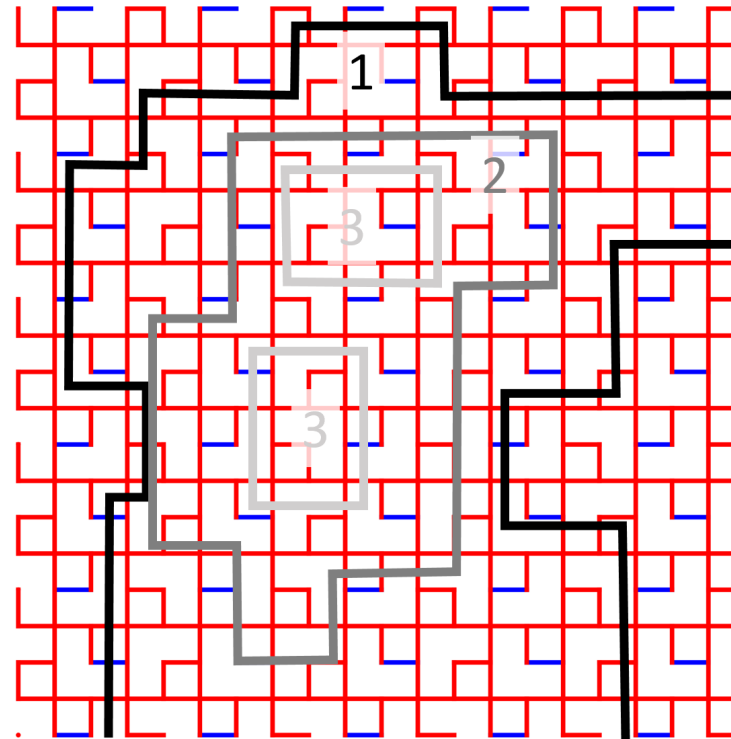
-Trees

-Fermions

-The Loops

-Re-weighting

Conclusion



The Loops

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

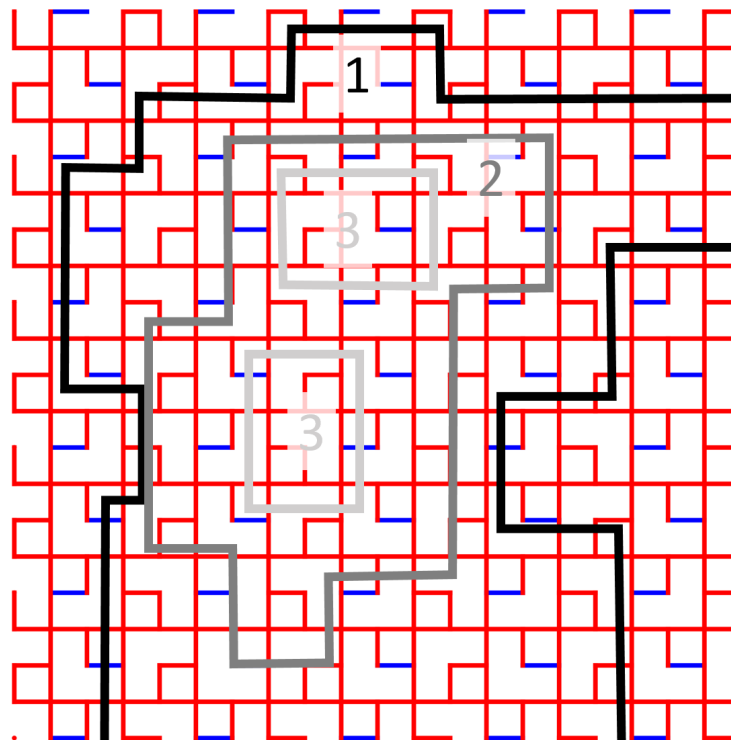
-Trees

-Fermions

-The Loops

-Re-weighting

Conclusion



- Each loop is weighted by two copies of signed restricted trees.

Re-weighting the Trees

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

-Fermions

-The Loops

-Re-weighting

Conclusion

- The partition function of the Free fermions is $Z = \det^2(\Delta)$.

Re-weighting the Trees

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

-Fermions

-The Loops

-Re-weighting

Conclusion

- The partition function of the Free fermions is $Z = \det^2(\Delta)$.
- We have to re-weigh the loops (in Girsanov sense) to account of

Re-weighting the Trees

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

-Fermions

-The Loops

-Re-weighting

Conclusion

- The partition function of the Free fermions is $Z = \det^2(\Delta)$.
- We have to re-weigh the loops (in Girsanov sense) to account of
 1. Sign.

Re-weighting the Trees

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

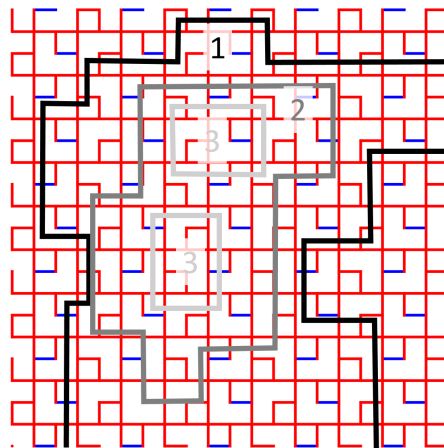
-Fermions

-The Loops

-Re-weighting

Conclusion

- The partition function of the Free fermions is $Z = \det^2(\Delta)$.
- We have to re-weigh the loops (in Girsanov sense) to account of
 1. Sign.
 2. Constraint



Re-weighting the Trees

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

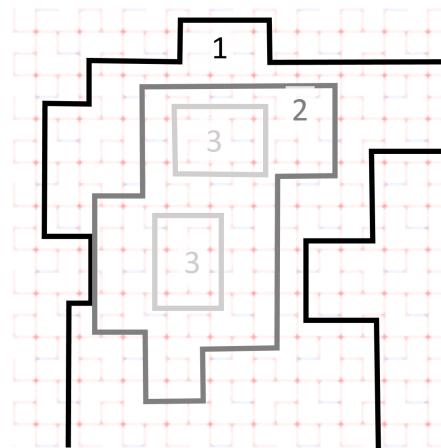
-Fermions

-The Loops

-Re-weighting

Conclusion

- The partition function of the Free fermions is $Z = \det^2(\Delta)$.
- We have to re-weigh the loops (in Girsanov sense) to account of
 1. Sign.
 2. Constraint



Re-weighting the Trees

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

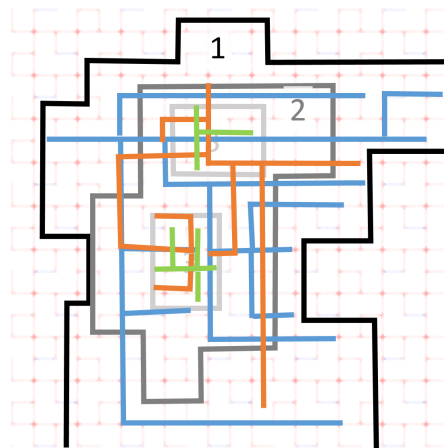
-Fermions

-The Loops

-Re-weighting

Conclusion

- The partition function of the Free fermions is $Z = \det^2(\Delta)$.
- We have to re-weigh the loops (in Girsanov sense) to account of
 1. Sign.
 2. Constraint



Re-weighting the Trees

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

-Fermions

-The Loops

-Re-weighting

Conclusion

- The partition function of the Free fermions is $Z = \det^2(\Delta)$.
- We have to re-weigh the loops (in Girsanov sense) to account of
 1. Sign.
 2. Constraint
- Both can be taken into account by bosonizing (introducing a height function that describes the trees). The height function is a Gaussian Free Field (GFF).

Re-weighting the Trees

Acknowledgement

Introduction

CLE Approach

-CLEs

-Manhattan

-Trees

-Fermions

-The Loops

-Re-weighting

Conclusion

- The partition function of the Free fermions is $Z = \det^2(\Delta)$.
- We have to re-weigh the loops (in Girsanov sense) to account of
 1. Sign.
 2. Constraint
- Both can be taken into account by bosonizing (introducing a height function that describes the trees). The height function is a Gaussian Free Field (GFF).
- $Z = \det^2(\Delta) \langle e^{\varphi A \varphi} \delta(\partial \varphi - v) \rangle_{GFF} \stackrel{?}{=} \det^{2+\delta c}(\Delta)$

Acknowledgement

Introduction

CLE Approach

Conclusion

Conclusion

Conclusion

Acknowledgement

Introduction

CLE Approach

Conclusion

- The integer quantum Hall transition may be described in the language of classical statistical mechanics.

Conclusion

Acknowledgement

Introduction

CLE Approach

Conclusion

- The integer quantum Hall transition may be described in the language of classical statistical mechanics.
- A loop ensemble may be developed.

Conclusion

Acknowledgement

Introduction

CLE Approach

Conclusion

- The integer quantum Hall transition may be described in the language of classical statistical mechanics.
- A loop ensemble may be developed.
- Final result for the weight of the loops is still pending.

Conclusion

Acknowledgement

Introduction

CLE Approach

Conclusion

- The integer quantum Hall transition may be described in the language of classical statistical mechanics.
- A loop ensemble may be developed.
- Final result for the weight of the loops is still pending.
- Critical exponents may be obtained as directly related to fractal dimensions of subsets of the loops and trees.