

# A Gutzwiller type trace formula for the magnetic Dirac operator

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Workshop on Geometric Quantization

# Coupled/Magnetic Dirac Operator

$(Y^{2n+1}, g^{TY})$  oriented, spin.

$(L, h^L)$  Hermitian line bundle with unitary connection  $A_0$ .

$a \in \Omega^1(Y)$  gives family of connections  $A_r = A_0 + ira$ ,  $r \in \mathbb{R}$ .

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Equivalently study  $D_h = hD_{A_0} + ic(a)$  as  $h = \frac{1}{r} \rightarrow 0$ .

# Quantization analog

$X^{2n}$  complex, Hermitian.  $L \rightarrow X$  Hermitian holomorphic line bundle

$$D_p = \sqrt{2} (\bar{\partial}_{L^p} + \bar{\partial}_{L^p}^*) : \Omega^{0,*}(L^p) \rightarrow \Omega^{0,*}(L^p).$$

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(Bordemann-Meinrenken  
-Schlichenmaier '94)

$$\begin{aligned} \|T_{f,p}\| &:= \|\Pi_p f\| \sim \|f\|_\infty, \\ [T_f, T_g] &= \frac{i}{p} T_{\{f,g\}} + O(p^{-1}) \end{aligned}$$

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Renormalized Laplacian:

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$\text{Spec}(\Delta_p) \subset [-M, M] \cup [C_1 p - C_2]; N_p = \underbrace{\#\text{Spec}(\Delta_p) \cap (-M, M)}_{\sim \text{vol}(X)p^n}$

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Proved using rescaling/local index theory arguments (cf. Ma-Marinescu '07).

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$$(Taubes '09) \eta(D_{A_r}) = o\left(r^{n+\frac{1}{2}+\varepsilon}\right). \quad (\text{Tsai '16}) \text{ in 3D}$$

Used in proof of 3D Weinstein conjecture via Seiberg-Witten equations.

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More generally;

$Y^{2m+1}$  = unit circle bundle of positive  $\mathcal{L}$



$X^{2m}$  (cpx)

Chern connection gives  $TY = TS^1 \oplus \pi^*TX$ ; choose adiabatic metrics

$$g_\varepsilon^{TY} = g^{TS^1} \oplus \frac{g^{TX}}{\varepsilon}, \quad \varepsilon > 0.$$

$L = \mathbb{C}$  with  $a = e^*$  = dual to  $S^1$  generator (contact with periodic Reeb flow)

# Circle bundle

Fourier decompr. along  $S^1$  writes  $\text{Spec}(D_{A_r})$  in terms of Dolbeault Laplacian  
 $\Delta_p^q : \Omega^{0,q}(X, \mathcal{L}^p) \rightarrow \Omega^{0,q}(X, \mathcal{L}^p)$ .

**Type 1.**

$$\lambda_r = (-1)^q \left[ p + \varepsilon \left( q - \frac{m}{2} \right) - r \right]$$

multiplicity =  $\dim H^q(X, \mathcal{K} \otimes \mathcal{L}^p)$ .

**Type 2.**

$$\lambda_r^\pm = \frac{(-1)^{q+1} \pm \sqrt{[2p + \varepsilon(2q - m) - r + 1]^2 + 4\mu^2\varepsilon}}{2}$$

$$0 < \frac{1}{2}\mu^2 \in \text{Spec}(\Delta_k^p).$$

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For  $\varepsilon \ll 1$  small can compute (Bismut-Cheeger '89, Dai '91)

$$\eta_r = \sum_{j=0}^m \left\{ \left( \frac{r^{j+1}}{(j+1)!} - \sum_{k=1}^{\lfloor r + \frac{\varepsilon m}{2} \rfloor} \frac{k^j}{j!} \right) \int_X c_1(\mathcal{L})^j [\text{ch}(\mathcal{K}) \text{td}(X)]^{m-j} \right\} + O(1)$$

Here  $\eta_r = O(r^n)$ . No leading term.

# Sharp results

To prove  $\eta_r = O(r^n)$

Set  $\mathfrak{J} \in C^\infty(i\mathfrak{u}(TX))$ , via  $g^{TY}(., \mathfrak{J}.) = da(., .)$

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## Theorem

(S '17) If further Reeb flow  $e^{tR}$  of  $a$  is non-resonant

$$\lim_{r \rightarrow \infty} r^{-n} \eta_r = -\frac{1}{2} \frac{1}{(2\pi)^{n+1}} \frac{1}{n!} \int_Y \left[ \text{tr } |\mathfrak{J}|^{-1} \left( \nabla^{TX} \mathfrak{J} \right) \right]$$

where  $|\mathfrak{J}| = \sqrt{-\mathfrak{J}^2}$  (restricted to  $R^\perp$ ).

Dynamical assumption excludes circle bundle example.

# Trace and $\eta$

Semi-classically:  $\eta_r = \eta_h = \eta(hD_{A_0} + c(a))$ ,  $h = \frac{1}{r} \rightarrow 0$ .

Using Bismut-Freed

$$\begin{aligned}\eta_h &= \int_0^\infty \frac{1}{\sqrt{\pi th}} \text{tr} \left( D_h e^{-\frac{t}{h} D_h^2} \right) dt \\ &= \int_0^1 \frac{1}{\sqrt{\pi th}} \text{tr} \left( D_h e^{-\frac{t}{h} D_h^2} \right) dt + \text{tr } E \left( \frac{1}{\sqrt{h}} D_h \right)\end{aligned}$$

Second term  $E(x) = \text{sign}(x) \operatorname{erfc}(|x|)$  has non-local trace.

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By Tauberian argument  $\eta_h$  asymptotics follow from

## Theorem

(Gutzwiller trace) For  $f, \theta \in C_c^\infty(\mathbb{R})$ ;  $\lambda \in \mathbb{R}$ ,

$$\begin{aligned}\text{tr} \left[ f \left( \frac{D_h}{\sqrt{h}} \right) \check{\theta} \left( \frac{\lambda \sqrt{h} - D_h}{h} \right) \right] &\sim h^{-m} \left( \sum_{j=0}^{\infty} c_j(\lambda) h^{j/2} \right) \\ &+ \sum_{\gamma} e^{\frac{i}{h} T_\gamma} e^{i \frac{\pi}{2} \mathfrak{m}_\gamma} \sum_{j=1}^{\infty} h^{j/2} \sum_{k=0}^j \lambda^k A_{\gamma,j,k} \theta(L_\gamma)\end{aligned}$$

where  $T_\gamma = \text{period}$ ,  $L_\gamma = \text{length}$ ,  $\mathfrak{m}_\gamma = \text{Maslov index of Reeb orbit}$ .



# Trace formula

For non-scalar operators trace formula known for diagonalizable symbol  $\sigma(x, \xi)$ .  
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Here  $\sigma(D_h) = c(\xi + a)$  with eigenvalues  $\pm |\xi + a|$ .  
Characteristic variety:  $\Sigma = \{\xi = -a\}$ .

# Normal form

Proof uses microlocal normal forms.

Contact 3D model:

$$D_0 = \begin{bmatrix} i & \\ & -i \end{bmatrix} h\partial_z + \underbrace{\begin{bmatrix} & -1 \\ 1 & \end{bmatrix} (h\partial_x + iy) + \begin{bmatrix} 0 & i \\ i & 0 \end{bmatrix} (h\partial_y - ix)}_{D_{00}}.$$

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More generally

$$D_h = \gamma^j w_j^k (h\partial_k + ia_k) + O(h)$$

where  $g_j^i = w_j^k w_k^i$  (diagonalizes metric).

$$\underbrace{UDU^*}_{\text{FIO conjugation}} = \begin{bmatrix} i & \\ & -i \end{bmatrix} h\partial_z + D_{00} + \underbrace{R}_{\text{commutes with } D_{00}^2} + O(h^\infty).$$

Study trace on the level.

# Koszul differentials

$U = e^{\frac{it}{\hbar} f^W} e^{it c_0^W(b)}$  with

$f^W$  = scalar Weyl-quantization;  $c_0^W(b)$  = Clifford-Weyl quantization of  
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Conjugation

$$e^{\frac{i}{\hbar} f} D_0 e^{-\frac{i}{\hbar} f} = D_0 + c_0 \left( \underbrace{\tilde{w}_\partial}_{\text{Koszul}} f \right) + \dots$$

$$e^{ic_0(b)} D_0 e^{-ic_0(b)} = D_0 + (-1)^{|a|} 2c_0 \left( \underbrace{i_x}_{\text{Koszul}} b \right) + hc_0 \left( \underbrace{\tilde{w}_\partial}_{\text{Koszul}} b \right) + \dots$$

$$i_x = xi_{e_1} + \xi i_{e_2}; \quad \tilde{w}_\partial = \partial_x e_2 \wedge -\partial_\xi e_1 \wedge$$

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Large time trace formula: similar normal form near (lifts to  $\Sigma$  of) Reeb orbits  
+ weak propagation lemma (Dimassi-Sjöstrand).

# Further questions

Still to do:

1. Relax assumptions (contact, on  $\text{Spec } \tilde{\mathcal{J}} \dots$ )
2. Asymptotics of  $\log \det D_h^2$  (expansion of low lying eigenvalues)
3. Understand propagator  $e^{it\frac{D}{\hbar}}$  & propagation of singularities
4. Apply technique to complex/CR analogs
5. Study microlocal analog (magnetic Laplacian  $\rightarrow$  subRiemannian Laplacian)

Thank you.