

Lattices and Quadratic Forms (Summer 2024) - Solutions to Problem Set 5

1. Given representatives $m \in L/cL$, we can uniquely write any element $n \in L + \mu$ as $n = r + m + \mu$ where $r \in cL$. Therefore,

$$\Theta_{L,\mu}(\tau/c) = \sum_{n \in L+\mu} e^{\pi i \frac{\tau}{c} n^2} = \sum_{m \in L/cL} \sum_{r \in cL} e^{\pi i \frac{\tau}{c} (r+m+\mu)^2} = \sum_{m \in L/cL} \sum_{r \in \sqrt{c}L} e^{\pi i \frac{\tau}{c} (\sqrt{c}r+m+\mu)^2}$$

and hence

$$\Theta_{L,\mu}(\tau/c) = \sum_{m \in L/cL} \sum_{r \in \sqrt{c}L + \frac{m+\mu}{\sqrt{c}}} e^{\pi i \tau r^2} = \sum_{m \in L/cL} \Theta_{\sqrt{c}L, \frac{1}{\sqrt{c}}(m+\mu)}(\tau),$$

proving the asserted identity. Here we also note that $\sqrt{c}L$ is a positive-definite, integral lattice and $\frac{1}{\sqrt{c}}(m+\mu) \in (\sqrt{c}L)^\sharp$. Further note that the summands $\Theta_{\sqrt{c}L, \frac{1}{\sqrt{c}}(m+\mu)}(\tau)$ are independent of the representatives chosen from L/cL elements and this is why we can view the sum as one over L/cL .

2. We start with the $(\tau, z) \mapsto (\tau+1, z)$ transformation. For θ_{01} , we have

$$\theta_{01}(\tau+1, z) = \sum_{m \in \mathbb{Z}} (-1)^m e^{\pi i m^2} e^{\pi i \tau m^2 + 2\pi i z m}.$$

Since $(-1)^m e^{\pi i m^2} = (-1)^{m(m+1)} = 1$, we find

$$\boxed{\theta_{01}(\tau+1, z) = \theta_{00}(\tau, z).}$$

For θ_{10} and θ_{11} , the transformation produces the $e^{\pi i(m+\frac{1}{2})^2}$ factor within the summands. Since

$$e^{\pi i(m+\frac{1}{2})^2} = e^{\pi i/4} (-1)^{m^2+m} = e^{\pi i/4} \quad \text{for } m \in \mathbb{Z},$$

we find

$$\boxed{\theta_{10}(\tau+1, z) = e^{\pi i/4} \theta_{10}(\tau, z)} \quad \text{and} \quad \boxed{\theta_{11}(\tau+1, z) = e^{\pi i/4} \theta_{11}(\tau, z).}$$

For the $(\tau, z) \mapsto (-\frac{1}{\tau}, \frac{z}{\tau})$ transformation, we start with the general expression

$$\theta_{ab}(\tau, z) = \sum_{m \in \mathbb{Z} + \frac{a}{2}} e^{\pi i \tau m^2 + 2\pi i m(z + \frac{b}{2})} = e^{-\pi i(z+b/2)^2/\tau} \sum_{m \in \mathbb{Z} + \frac{a}{2}} e^{\pi i \tau(m + \frac{z+b/2}{\tau})^2}$$

so that

$$\theta_{ab}\left(-\frac{1}{\tau}, \frac{z}{\tau}\right) = e^{\pi i(z+b\tau/2)^2/\tau} \sum_{m \in \mathbb{Z} + \frac{a}{2}} e^{-\frac{\pi i}{\tau}(m - z - \frac{b\tau}{2})^2} = e^{\pi i(z+b\tau/2)^2/\tau} \sum_{m \in \mathbb{Z}} f_\tau\left(m - z - \frac{b\tau - a}{2}\right),$$

where

$$f_\tau(x) := e^{-\frac{\pi i}{\tau} x^2}.$$

This Gaussian function is a Schwartz function and we use Poisson summation (on the self-dual lattice \mathbb{Z}) while noting that the Fourier transform¹ of $f_\tau(x-r)$ is $\sqrt{-i\tau} e^{\pi i \tau k^2 - 2\pi i k r}$ to find²

$$\theta_{ab}\left(-\frac{1}{\tau}, \frac{z}{\tau}\right) = e^{\pi i(z+b\tau/2)^2/\tau} \sqrt{-i\tau} \sum_{k \in \mathbb{Z}} e^{\pi i \tau k^2 - 2\pi i k(z + \frac{b\tau - a}{2})}.$$

¹The Fourier transform $\hat{f}(k)$ of $f(x)$ is

$$\hat{f}(k) = \int_{\mathbb{R}} dx f(x) e^{-2\pi i k x}.$$

²This is including the case where r is complex as can be confirmed by performing the relevant Gaussian integrals.

Taking $k \mapsto -k$, we rewrite this as

$$\theta_{ab}\left(-\frac{1}{\tau}, \frac{z}{\tau}\right) = e^{\pi i(z+b\tau/2)^2/\tau} \sqrt{-i\tau} \sum_{k \in \mathbb{Z}} e^{\pi i \tau k^2 + 2\pi i k(z + \frac{b\tau-a}{2})}$$

and then note

$$e^{\frac{\pi i}{\tau}(z+b\tau/2)^2} e^{\pi i \tau k^2 + 2\pi i k(z + \frac{b\tau-a}{2})} = e^{\frac{\pi i}{\tau} z^2} e^{\pi i \tau (k + \frac{b}{2})^2} e^{2\pi i (k + \frac{b}{2})(z + \frac{a}{2})} e^{-2\pi i k a - \frac{\pi i}{2} a b}.$$

Since $a \in \{0, 1\}$, we have $e^{-2\pi i k a} = 1$ for $k \in \mathbb{Z}$ and this leads to

$$\theta_{ab}\left(-\frac{1}{\tau}, \frac{z}{\tau}\right) = e^{-\frac{\pi i}{2} a b} e^{\pi i z^2/\tau} \sqrt{-i\tau} \sum_{k \in \mathbb{Z} + \frac{b}{2}} e^{\pi i \tau k^2 + 2\pi i k(z + \frac{a}{2})}$$

and hence

$$\theta_{ab}\left(-\frac{1}{\tau}, \frac{z}{\tau}\right) = e^{-\frac{\pi i}{2} a b} e^{\pi i z^2/\tau} \sqrt{-i\tau} \theta_{ba}(\tau, z).$$

In particular,

$$\theta_{01}\left(-\frac{1}{\tau}, \frac{z}{\tau}\right) = e^{\pi i z^2/\tau} \sqrt{-i\tau} \theta_{10}(\tau, z), \quad \theta_{10}\left(-\frac{1}{\tau}, \frac{z}{\tau}\right) = e^{\pi i z^2/\tau} \sqrt{-i\tau} \theta_{01}(\tau, z),$$

and

$$\theta_{11}\left(-\frac{1}{\tau}, \frac{z}{\tau}\right) = -i e^{\pi i z^2/\tau} \sqrt{-i\tau} \theta_{11}(\tau, z).$$

3. Under $\tau \mapsto \tau + 1$ we have

$$\Psi_{\boldsymbol{\mu}}(\tau + 1) = \sum_{\mathbf{n} \in L + \boldsymbol{\mu} + \frac{\boldsymbol{\ell}}{2}} e^{\pi i \mathbf{n}^2} e^{\pi i \tau \mathbf{n}^2 + \pi i \mathbf{n} \cdot \boldsymbol{\ell}}.$$

Now note that for $\mathbf{n} = \mathbf{m} + \boldsymbol{\mu} + \frac{\boldsymbol{\ell}}{2}$ with $\mathbf{m} \in L$ we have

$$e^{\pi i (\mathbf{m} + \boldsymbol{\mu} + \frac{\boldsymbol{\ell}}{2})^2} = e^{\pi i (\boldsymbol{\mu} + \frac{\boldsymbol{\ell}}{2})^2 + \pi i (\mathbf{m}^2 + \mathbf{m} \cdot \boldsymbol{\ell}) + 2\pi i \mathbf{m} \cdot \boldsymbol{\mu}}.$$

Since $\boldsymbol{\mu} \in L^\sharp$ we have $\mathbf{m} \cdot \boldsymbol{\mu} \in \mathbb{Z}$ and since $\boldsymbol{\ell}$ is a characteristic vector in the integral lattice L we have $\mathbf{m}^2 + \mathbf{m} \cdot \boldsymbol{\ell} \in 2\mathbb{Z}$. Therefore, $e^{\pi i (\mathbf{m} + \boldsymbol{\mu} + \frac{\boldsymbol{\ell}}{2})^2} = e^{\pi i (\boldsymbol{\mu} + \frac{\boldsymbol{\ell}}{2})^2}$ factors out as an overall factor and we get

$$\Psi_{\boldsymbol{\mu}}(\tau + 1) = e^{\pi i (\boldsymbol{\mu} + \frac{\boldsymbol{\ell}}{2})^2} \Psi_{\boldsymbol{\mu}}(\tau).$$

To study the behavior under the inversion $\tau \mapsto -\frac{1}{\tau}$ we first rewrite

$$\Psi_{\boldsymbol{\mu}}(\tau) = \sum_{\mathbf{n} \in L + \boldsymbol{\mu} + \frac{\boldsymbol{\ell}}{2}} e^{\pi i \tau \mathbf{n}^2 + \pi i \mathbf{n} \cdot \boldsymbol{\ell}} = e^{-\pi i \frac{\boldsymbol{\ell}^2}{4\tau}} \sum_{\mathbf{n} \in L + \boldsymbol{\mu} + \frac{\boldsymbol{\ell}}{2}} e^{\pi i \tau (\mathbf{n} + \frac{\boldsymbol{\ell}}{2\tau})^2} = e^{-\pi i \frac{\boldsymbol{\ell}^2}{4\tau}} \sum_{\mathbf{n} \in L} e^{\pi i \tau (\mathbf{n} + \boldsymbol{\mu} + \frac{\boldsymbol{\ell}}{2} + \frac{\boldsymbol{\ell}}{2\tau})^2}$$

so that

$$\Psi_{\boldsymbol{\mu}}\left(-\frac{1}{\tau}\right) = e^{\pi i \tau \frac{\boldsymbol{\ell}^2}{4}} \sum_{\mathbf{n} \in L} f_{\tau}\left(\mathbf{n} + \boldsymbol{\mu} + \frac{\boldsymbol{\ell}}{2} - \frac{\boldsymbol{\ell}\tau}{2}\right)$$

where

$$f_{\tau}(\mathbf{x}) := e^{-\frac{\pi i}{\tau} \mathbf{x}^2}.$$

This Gaussian function is a Schwartz function and we use Poisson summation while noting that the Fourier transform of $f_{\tau}(\mathbf{x} + \mathbf{r})$ is $(-i\tau)^{n/2} e^{\pi i \tau \mathbf{k}^2 + 2\pi i \mathbf{k} \cdot \mathbf{r}}$ (also when \mathbf{r} is complex) to find

$$\Psi_{\boldsymbol{\mu}}\left(-\frac{1}{\tau}\right) = e^{\pi i \tau \frac{\boldsymbol{\ell}^2}{4}} \frac{(-i\tau)^{n/2}}{\sqrt{\det L}} \sum_{\mathbf{k} \in L^\sharp} e^{\pi i \tau \mathbf{k}^2 + 2\pi i \mathbf{k} \cdot (\boldsymbol{\mu} + \frac{\boldsymbol{\ell}}{2} - \frac{\boldsymbol{\ell}\tau}{2})}.$$

Letting $\mathbf{k} \mapsto -\mathbf{k}$ and then writing the sum $\sum_{\mathbf{k} \in L^\sharp}$ as $\sum_{\boldsymbol{\nu} \in L^\sharp/L} \sum_{\mathbf{k} \in L+\boldsymbol{\nu}}$ we find

$$\Psi_\mu\left(-\frac{1}{\tau}\right) = e^{\pi i \tau \frac{\ell^2}{4}} \frac{(-i\tau)^{n/2}}{\sqrt{\det L}} \sum_{\boldsymbol{\nu} \in L^\sharp/L} \sum_{\mathbf{k} \in L+\boldsymbol{\nu}} e^{\pi i \tau \mathbf{k}^2 - 2\pi i \mathbf{k} \cdot (\mu + \frac{\ell}{2} - \frac{\ell\tau}{2})}.$$

Now note that

$$e^{\pi i \tau \frac{\ell^2}{4}} e^{\pi i \tau \mathbf{k}^2 - 2\pi i \mathbf{k} \cdot (\mu + \frac{\ell}{2} - \frac{\ell\tau}{2})} = e^{\pi i \tau (\mathbf{k} + \frac{\ell}{2})^2} e^{-\pi i \mathbf{k} \cdot \ell} e^{-2\pi i \mathbf{k} \cdot \mu},$$

where

$$e^{-2\pi i \mathbf{k} \cdot \mu} = e^{-2\pi i \boldsymbol{\nu} \cdot \mu} \quad \text{for } \mathbf{k} \in L + \boldsymbol{\nu}.$$

Moreover, for any $\mathbf{k} \in L + \boldsymbol{\nu}$, which we can write as $\mathbf{k} = \mathbf{m} + \boldsymbol{\nu}$ with $\mathbf{m} \in L$, we have $\mathbf{k} \cdot \ell = (\mathbf{m} + \boldsymbol{\nu}) \cdot \ell \in \mathbb{Z}$ since $\mathbf{m} \cdot \ell \in \mathbb{Z}$ for $\mathbf{m}, \ell \in L$ with L an integral lattice and since $\boldsymbol{\nu} \cdot \ell \in \mathbb{Z}$ because $\boldsymbol{\nu}$ is in the dual lattice. Therefore,

$$e^{-\pi i \mathbf{k} \cdot \ell} = e^{+\pi i \mathbf{k} \cdot \ell} = e^{\pi i (\mathbf{k} + \frac{\ell}{2}) \cdot \ell - \frac{\pi i}{2} \ell^2}. \quad (0.1)$$

Consequently, we find

$$\Psi_\mu\left(-\frac{1}{\tau}\right) = e^{-\frac{\pi i}{2} \ell^2} \frac{(-i\tau)^{n/2}}{\sqrt{\det L}} \sum_{\boldsymbol{\nu} \in L^\sharp/L} e^{-2\pi i \boldsymbol{\mu} \cdot \boldsymbol{\nu}} \sum_{\mathbf{k} \in L+\boldsymbol{\nu}+\frac{\ell}{2}} e^{\pi i \tau \mathbf{k}^2 + \pi i \mathbf{k} \cdot \ell}$$

and hence

$$\boxed{\Psi_\mu\left(-\frac{1}{\tau}\right) = e^{-\frac{\pi i}{2} \ell^2} \frac{(-i\tau)^{n/2}}{\sqrt{\det L}} \sum_{\boldsymbol{\nu} \in L^\sharp/L} e^{-2\pi i \boldsymbol{\mu} \cdot \boldsymbol{\nu}} \Psi_\nu(\tau).}$$

4. Let us consider the one dimensional integral lattice $L \simeq \mathbb{Z}$ with the symmetric bilinear form $B(n, m) = 3nm$ for $n, m \in \mathbb{Z}$. We have $\det L = 3$ with representatives of the discriminant group L^\sharp/L given by $\mu = 0, +\frac{1}{3}, -\frac{1}{3}$. Also noting that $\ell = 1$ is a characteristic vector for this lattice (since $3n^2 + 3n \in 2\mathbb{Z}$), the corresponding theta functions $\Psi_\mu(\tau)$ (of the form we considered in Problem 3) are

$$\begin{aligned} \Psi_0(\tau) &= \sum_{n \in \mathbb{Z}} e^{3\pi i \tau (n + \frac{1}{2})^2 + 3\pi i (n + \frac{1}{2})} = -i \sum_{n \in \mathbb{Z}} (-1)^n e^{3\pi i \tau (n + \frac{1}{2})^2} \\ \Psi_{1/3}(\tau) &= \sum_{n \in \mathbb{Z}} e^{3\pi i \tau (n + \frac{5}{6})^2 + 3\pi i (n + \frac{5}{6})} = i \sum_{n \in \mathbb{Z}} (-1)^n e^{3\pi i \tau (n + \frac{5}{6})^2} \\ \Psi_{-1/3}(\tau) &= \sum_{n \in \mathbb{Z}} e^{3\pi i \tau (n + \frac{1}{6})^2 + 3\pi i (n + \frac{1}{6})} = i \sum_{n \in \mathbb{Z}} (-1)^n e^{3\pi i \tau (n + \frac{1}{6})^2}. \end{aligned}$$

Importantly, we note that

$$\Psi_{-1/3}(\tau) = i \eta(\tau).$$

Moreover, by letting $n \mapsto -n - 1$ we find

$$\Psi_0(\tau) = 0 \quad \text{and} \quad \Psi_{1/3}(\tau) = -i \eta(\tau).$$

Therefore, the corresponding transformation in Problem 3 implies

$$\Psi_{-1/3}(-1/\tau) = e^{-\frac{3\pi i}{2}} \frac{(-i\tau)^{1/2}}{\sqrt{3}} \left(\Psi_0(\tau) + e^{-6\pi i (-\frac{1}{3})(-\frac{1}{3})} \Psi_{-1/3}(\tau) + e^{-6\pi i (-\frac{1}{3})(\frac{1}{3})} \Psi_{1/3}(\tau) \right).$$

This yields

$$i \eta(-1/\tau) = i \frac{(-i\tau)^{1/2}}{\sqrt{3}} \left(i e^{-\frac{2\pi i}{3}} \eta(\tau) - i e^{\frac{2\pi i}{3}} \eta(\tau) \right).$$

Since $i e^{-\frac{2\pi i}{3}} - i e^{\frac{2\pi i}{3}} = \sqrt{3}$, we finally find

$$\boxed{\eta(-1/\tau) = \sqrt{-i\tau} \eta(\tau)}$$

as claimed.